



# Nonlinear collision effect on alpha particle confinement

Yoshitada Masaoka and Sadayoshi Murakami

Department of Nuclear Engineering Kyoto University

## Abstract

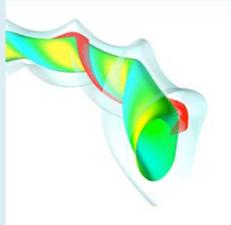
Confinement of  $\alpha$ -particles is investigated including the collisions with various plasma species such as electron, deuterium, tritium, and high-energy  $\alpha$ -particle itself in a heliotron fusion reactor, which is based on the LHD configurations. GNET (Global Neoclassical Transport) code is being improved to take into account the nonlinear collision effect on the  $\alpha$ -particle confinement. The code is benchmarking with the linear operator in the shifted Maxwellian plasma.

## Introduction

### Helical device

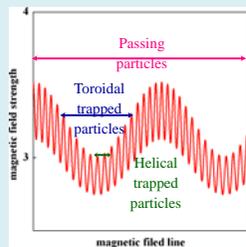
- The magnetic field is generated mainly by the coil current.
  - ◆ Permits a steady state plasma.
  - ◆ No plasma disruption caused by the plasma current.
- The magnetic configuration is inherently three-dimensional (3D).
  - ◆ The plasma behavior is more complex than in tokamaks.

† Several physics and technical problems remain to be studied and solved, such as the behavior and confinement of high energy  $\alpha$  particles in helical plasma.



### $\alpha$ particle in helical plasma

- Helical trapped particles : trapped in the helical ripple
- Toroidal trapped particles : trapped in the toroidal ripple
- Passing particles : not trapped in either the helical or toroidal ripples
- Transition particles : transition between being trapped particles and passing particles



➔ These trapped motions cause complex orbits of trapped particles and enhance radial diffusion of energetic particles.

## Simulation model

### GNET code

We solve **the drift kinetic equation** in the 5D phase-space with pitch angle and energy scattering using the **GNET code (Global Neoclassical Transport code)** [4].

### The drift kinetic equation

$$\frac{\partial f_\alpha}{\partial t} + (\mathbf{v}_\parallel + \mathbf{v}_D) \cdot \nabla f_\alpha + \dot{\mathbf{v}} \cdot \nabla_{\mathbf{v}} f_\alpha = C^{\text{coll}}(f_\alpha) + L^{\text{particle}}(f_\alpha) + S_\alpha$$

- $f_\alpha$  : distribution function of  $\alpha$  particles
- $v_\parallel$  : velocity parallel to magnetic field line
- $v_D$  : drift velocity
- $C^{\text{coll}}$  : Coulomb collision operator (linear and nonlinear)
- $L^{\text{particle}}$  : particle loss term (LCFS)
- $S_\alpha$  : particle source generated by fusion reaction

The steady state distribution of  $\alpha$  particle is evaluated. The GNET code uses a Monte Carlo technique to calculate the distribution function of a set of test particles.

### $\alpha$ particle source ( $S_\alpha$ )

#### Fusion reaction rate

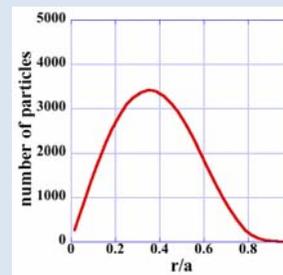
$$S_\alpha = n_D n_T \langle \sigma v \rangle = n_D n_T \int \int f_D(v_D) f_T(v_T) \sigma_T(E) |v_D - v_T| dv_D dv_T$$

- $\sigma$  : total reaction cross-section
- $E$  : relative energy
- $n_e$  : radial profile of plasma density
- $T_e$  : plasma temperature

$$n_e(\rho) [10^{20} \text{m}^{-3}] = 1.9(1 - \rho^8) + 0.1$$

$$T_e(\rho) [\text{keV}] = 9.5(1 - \rho^2) + 0.5$$

- $\rho$  : normalized minor radius
- $0$  : the value at the magnetic axis
- $1$  : the value at the last closed flux surface (LCFS)



Based on the fusion reaction rate, we get **an initial radial profile of  $\alpha$  particles.**

## The nonlinear collision operator : $C_\alpha^{\text{nonlinear}}$

We can write the  $C_\alpha^{\text{nonlinear}}$  with Rosenbluth potentials[6],

$$C^{a/b} = -D \overset{\leftrightarrow}{\nabla} f_a(\mathbf{v}) + F \overset{\leftrightarrow}{\nabla} f_a(\mathbf{v})$$

$$D = -\frac{4\pi\Gamma^{a/b}}{n_b} \nabla \nabla \psi_b(\mathbf{v})$$

$$F = -\frac{4\pi\Gamma^{a/b}}{n_b} \frac{m_a}{m_b} \nabla \phi_b(\mathbf{v})$$

$$\Gamma^{a/b} = \frac{n_b q_a^2 q_b^2 \ln \Lambda^{a/b}}{4\pi \epsilon_0^2 m_a^2}$$

- $D$  : diffusion tensor
- $F$  : average force tensor
- $n$  : density of plasma
- $m$  : mass of particle species
- $\epsilon_0$  : electrical constant
- $a$  : test particle species
- $b$  : background particle species

### ➤ Rosenbluth potentials $\phi, \psi$

$$\phi_a(v, \theta) = \sum_{l=0}^{\infty} \sum_b \frac{m_a + m_b}{m_b} \phi_b^{(l)}(v) P_l(\cos \theta)$$

$$\psi_a(v, \theta) = \sum_{l=0}^{\infty} \sum_b \psi_b^{(l)}(v) P_l(\cos \theta)$$

$$\phi_b^{(l)}(v) = -\frac{1}{2l+1} \left[ \int_0^v \frac{v'^{l+2}}{v'^{l+1}} f_b^{(l)}(v') dv' + \int_v^\infty \frac{v'^l}{v'^{l-1}} f_b^{(l)}(v') dv' \right]$$

$$\psi_b^{(l)}(v) = \frac{1}{2(4l^2 - 1)} \left[ \int_0^v \frac{v'^{l+2}}{v'^{l-1}} \left( 1 - \frac{(l - \frac{1}{2}) v'^2}{(l + \frac{3}{2}) v'^2} \right) f_b^{(l)}(v') dv' \right. \\ \left. + \int_v^\infty \frac{v'^l}{v'^{l-3}} \left( 1 - \frac{(l - \frac{1}{2}) v'^2}{(l + \frac{3}{2}) v'^2} \right) f_b^{(l)}(v') dv' \right]$$

### ➤ Legendre Polynomial Expansion

$$f(v, \theta) = \sum_{l=0}^{\infty} f^{(l)}(v) P_l(\cos \theta)$$

$$f^{(l)}(v) = \frac{2l+1}{2} \int_0^\pi f(v, \theta) P_l(\cos \theta) \sin \theta d\theta$$

$$P_0(\mu) = 1$$

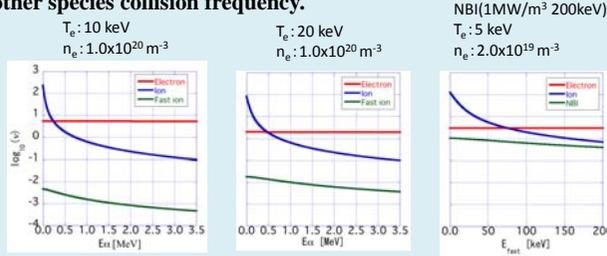
$$(l+1)P_{l+1}(\mu) = (2l+1)\mu P_l(\mu) - lP_{l-1}(\mu)$$

## Nonlinear collision effect

- The relative velocity between high-energy particles sometimes becomes very small.
- Although the amount of high-energy particles are much less than thermal ions, it is considered that the nonlinear collision by each fast ion has usually larger effect than that by other background ions [1].
- This collision effect may lead to deteriorate the high-energy particle confinement, because of increasing a pitch angle scattering.

## Collision frequency

We compare with the beam-beam collision frequency and the beam-other species collision frequency.



## Objective

• Assuming **LHD type reactor** as a typical helical reactor, we investigate the helical fusion reactor **in a view point of the  $\alpha$ -particle confinement**.

• We include the collisional effects (**the energy and pitch angle scattering**) and evaluated **the distribution function of  $\alpha$ -particles**.

• We analyze including the both complicated orbit and nonlinear collision effects in order to make clear the  $\alpha$ -particle confinement in heliotrons.

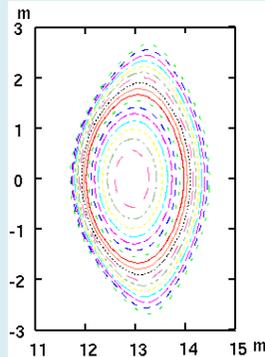
• The assumed fusion reactor

• The helical type of fusion reactor extending the LHD magnetic configuration. ( $R_{ax}$  is about 3.55 times larger than that of the LHD.)

Plasma volume : **1000m<sup>3</sup>**  
Magnetic field : **5T**

Magnetic configuration

( $R_{ax}$  : the magnetic axis position in vacuum): **NC**, which is **the neoclassical transport optimized configuration, based on  $R_{ax}=3.53m$  of LHD [3]**.



## Coulomb collision ( $C^{coll}(f)$ )

$C$  is the Coulomb collision operator including the linear collision effect  $C^{linear}$  and the nonlinear collision effect  $C^{nonlinear}$ .

$$C^{coll}(f_\alpha) = C_e^{linear}(f_\alpha) + C_D^{linear}(f_\alpha) + C_T^{linear}(f_\alpha) + C_\alpha^{linear}(f_\alpha) + C_\alpha^{nonlinear}(f_\alpha)$$

**The linear collision operator :  $C_i^{linear}$**

The operator of the pitch angle and energy scattering with background ions and electrons[5].

Pitch angle :

$$\lambda_n = \lambda_{n-1} - \sum_i \left( \nu_d^i \tau^i \mp [(1 - \lambda_{n-1}^2) \nu_d^i \tau^i]^{1/2} \right)$$

Energy :

$$E_n = E_{n-1} - \sum_i \left( (2\nu^i \tau^i) \left[ E_{n-1} \left( \frac{3}{2} + \frac{E_{n-1}}{\nu^i} \frac{d\nu^i}{dE_{n-1}} \right) \mp 2 \{ E_T E_{n-1} (\nu^i \tau^i) \}^{1/2} \right] \right)$$

$i$  : background ions(D, T,  $\alpha$ ) and electrons

$\nu_d$  : the deflection collision frequency

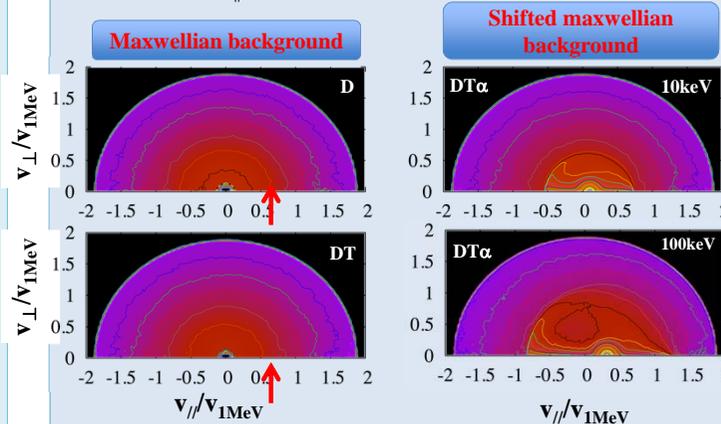
$\tau$  : the length of a time step

$n, n-1$  : numbers of time step

$\pm$  : the signs to be chosen randomly

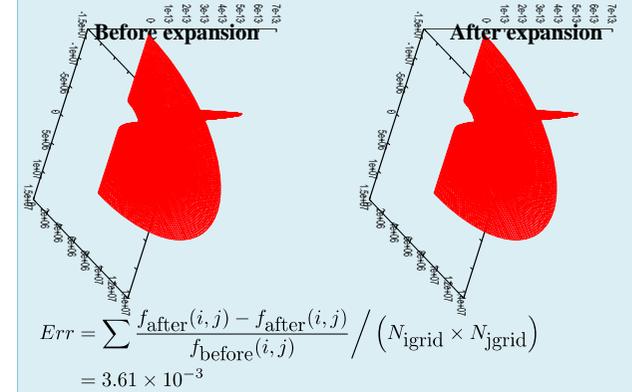
$E_n, E_T$  : the test particle energy at time step and thermal energy

$$\lambda = \nu_{||}/\nu$$



We evaluate the velocity space distribution of  $\alpha$  particle in the case with the linear collision operator (no orbit calculation).

- An increase of mass density of back ground plasma leads to improvement of slowing down.
- A shifted maxwellian background (10keV, 100keV) shifts the velocity space distributions.
- We will benchmark the nonlinear collision operator using the same background distributions.



Solving this collision operator at each time step, it take a lot of time in the simulation.

Therefore, before the simulation run, we build a database which give us the  $\alpha$ -particle velocity changes as a function of  $\alpha$ -particle velocity.

- We evaluate the Rosenbluth potential at the grid point in the velocity space.
- Using these Rosenbluth potentials, we build collision operator database.

## Summary

We improve the GNET code to take into account the nonlinear collision effect on the  $\alpha$ -particle confinement.

- We have extended the linear collision operator to estimate the effect of multi species plasma (deuterium, tritium, and alpha particle).
- We have studied the nonlinear collision operator and obtained its diffusion equation.
- The code is still need improvements.

## Reference

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